

ON THE 4×4 MATRICES ASSOCIATED WITH THE QUATERNIONIC UNITS

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ABSTRACT. Here we consider real and unitary quaternions which generate 3-rotations, and we show the corresponding 4×4 matrices associated with the quaternionic units.

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1. INTRODUCTION

For an arbitrary matrix $M_{4 \times 4}$ it is possible to obtain its expression in terms of the sixteen Dirac matrices in the standard representation [1, 2, 3, 4, 5]:

$$(1) \quad I, \quad \gamma^\mu, \quad \gamma^5, \quad \gamma^\mu \gamma^5, \quad \sigma^{\mu\nu},$$

in fact:

$$(2) \quad M = \begin{pmatrix} A & B \\ C & D \end{pmatrix} = e_0 \gamma^0 \gamma^5 + b_0 I + c_0 \gamma^0 + i d_0 \gamma^5 + b_1 \sigma^{23} + b_2 \sigma^{31} + b_3 \sigma^{12} + \\ + \sum_{j=1}^3 \left(e_j \gamma^j + c_j \gamma^j \gamma^5 + d_j \sigma^{0j} \right), \\ = \sum_{k=0}^3 \left[e_k \begin{pmatrix} 0 & \sigma_k \\ -\sigma_k & 0 \end{pmatrix} + b_k \begin{pmatrix} \sigma_k & 0 \\ 0 & \sigma_k \end{pmatrix} + c_k \begin{pmatrix} \sigma_k & 0 \\ 0_k & -\sigma_k \end{pmatrix} + i d_k \begin{pmatrix} 0 & \sigma_k \\ \sigma_k & 0 \end{pmatrix} \right],$$

in terms of the Pauli matrices with $\sigma_0 = I_{2 \times 2}$, therefore:

$$A = \sum_{k=0}^3 (b_k + c_k) \sigma_k, \quad B = \sum_{k=0}^3 (e_k + i d_k) \sigma_k, \\ C = \sum_{k=0}^3 (-e_k + i d_k) \sigma_k, \quad D = \sum_{k=0}^3 (b_k - c_k) \sigma_k,$$

whose solution is given by:

$$e_k = \frac{1}{4} \operatorname{tr} \left[(B - C) \sigma_k \right], \quad d_k = -\frac{i}{4} \operatorname{tr} \left[(B + C) \sigma_k \right],$$

$$k = 0, 1, 2, 3,$$

$$(3) \quad c_k = \frac{1}{4} \operatorname{tr} \left[(A - D) \sigma_k \right], \quad b_k = \frac{1}{4} \operatorname{tr} \left[(A + D) \sigma_k \right].$$

That is, if we have M , therefore we know A, B, C, D and we can calculate $(A \pm D)$ and $(B \pm C)$ and their products with the Pauli matrices, finally we obtain the corresponding traces to determine the coefficients of the expansion (2) in according with (3). For example, M could be any Lorentz matrix [6, 7], or the matrix that transforms a Dirac spinor under Lorentz mappings [8, 9], the Frenet–Serret curvature matrix for the motion of point particles in Minkowski spacetime [10, 11, 12], or the Faraday complex matrix [13], etc.

Here we consider real and unitary quaternions [14, 15, 16, 17]:

$$(4) \quad \mathbf{A} = a_0 + a_1 \mathbf{I} + a_2 \mathbf{J} + a_3 \mathbf{K}, \quad \bar{\mathbf{A}} = a_0 - a_1 \mathbf{I} - a_2 \mathbf{J} - a_3 \mathbf{K}, \quad a_0^2 + a_1^2 + a_2^2 + a_3^2 = 1,$$

with which we can associate the following 4×4 matrix [18, 19]:

$$(5) \quad M = \begin{pmatrix} a_0 & a_1 & a_2 & a_3 \\ -a_1 & a_0 & -a_3 & a_2 \\ -a_2 & a_3 & a_0 & -a_1 \\ -a_3 & -a_2 & a_1 & a_0 \end{pmatrix}$$

in terms of the Euler–Olinde Rodrigues parameters [20, 21, 22, 23]. In Sec. 2 we employ the relations (2), ..., (5) to determine the corresponding 4×4 matrices associated with the quaternionic units.

2. QUATERNIONIC UNITS AND ITS ASSOCIATION WITH 4×4 MATRICES

If we use (5) in (3) we obtain that the only non-zero values are:

$$(6) \quad b_0 = a_0, \quad c_2 = ia_1, \quad d_2 = a_3, \quad e_0 = a_2,$$

then (2) gives the expansion:

$$(7) \quad M = a_0 I + ia_1 \gamma^2 \gamma^5 + a_2 \gamma^0 \gamma^5 + a_3 \sigma^{02},$$

which implies the following matrix association for the quaternionic units:

$$(8) \quad 1 \rightarrow I, \quad \mathbf{I} \rightarrow i\gamma^2 \gamma^5, \quad \mathbf{J} \rightarrow \gamma^0 \gamma^5, \quad \mathbf{K} \rightarrow \sigma^{02},$$

in harmony with the expressions:

$$(9) \quad \mathbf{IJ} = \mathbf{K}, \quad \mathbf{JK} = \mathbf{I}, \quad \mathbf{KI} = \mathbf{J}, \quad \mathbf{I}^2 = \mathbf{J}^2 = \mathbf{K}^2 = -1.$$

On the other hand, the quaternion (4) allows generate arbitrary 3-rotations [24, 25] via the relation [19, 26, 27, 28, 29, 30, 31]:

$$(10) \quad \tilde{\mathbf{X}} = \mathbf{A} \mathbf{X} \bar{\mathbf{A}}, \quad \mathbf{X} = ct - ix\mathbf{I} - iy\mathbf{J} - iz\mathbf{K},$$

where all a_ν are real, then in (10) we can apply the matrix association (7) to obtain that $\tilde{t} = t$ and:

$$(11) \quad \begin{pmatrix} \tilde{x} \\ \tilde{y} \\ \tilde{z} \end{pmatrix} = R \begin{pmatrix} x \\ y \\ z \end{pmatrix}, \quad RR^T = I,$$

such that the matrix R represents a passive rotation (in the terminology of Ryder [32]) [19, 30, 31, 33, 34, 35]:

$$(12) \quad R = \begin{pmatrix} 1 - 2(a_2^2 + a_3^2) & 2(a_1a_2 - a_0a_3) & 2(a_1a_3 + a_0a_2) \\ 2(a_1a_2 + a_0a_3) & 1 - 2(a_1^2 + a_3^2) & 2(a_2a_3 - a_0a_1) \\ 2(a_1a_3 - a_0a_2) & 2(a_0a_1 + a_2a_3) & 1 - 2(a_1^2 + a_2^2) \end{pmatrix}.$$

Remark.- The expression (5) can be obtained from (7) and the relations [1, 2, 4, 5, 36, 37]:

$$(13) \quad \gamma^2\gamma^5 = \begin{pmatrix} \sigma_2 & 0 \\ 0 & -\sigma_2 \end{pmatrix}, \quad \gamma^0\gamma^5 = \begin{pmatrix} 0 & I \\ -I & 0 \end{pmatrix}, \quad \sigma^{02} = i \begin{pmatrix} 0 & \sigma_2 \\ \sigma_2 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}.$$

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